

**THE BMO^{-1} SPACE AND ITS APPLICATION TO
 SCHECHTER' S INEQUALITY**

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ABSTRACT. In this work, we study quadratic form inequalities of Schechter type; i.e., we characterize f for which there exists a positive constant C such that, for every $\epsilon \in (0, \infty)$,

$$\left| \int |u|^2 f dx \right| \leq \epsilon \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + C\epsilon^{-\beta} \|u\|_{L^2(\mathbb{R}^d)}^2, u \in C_0^\infty(\mathbb{R}^d), 0 < \beta < 1$$

Such quadratic form inequalities can be understood entirely in the framework of BMO^{-1} , using mean oscillations of $\nabla \Delta^{-1} f$ on balls. We show that this inequality holds if and only if $f \in BMO^{-1}(\mathbb{R}^d)$ if $\beta = 1$ or respectively if f lies in the homogeneous Besov space $\dot{B}_\infty^{-\frac{2\beta}{1+\beta}, \infty}$ if $0 < \beta < 1$.

1. INTRODUCTION

In this paper, we characterize the class of potentials $f \in \mathcal{D}'(\mathbb{R}^d)$ such that the quadratic form $\langle f, \cdot, \cdot \rangle$ has zero relative bound with respect to $H_0 = -\Delta$ on $L^2(\mathbb{R}^d)$ (see [8], X.17). In other words, for $f(x) \geq 0$ in $L^1_{loc}(\mathbb{R}^d)$, this property can be expressed in the form of the integral inequality :

$$(1.1) \quad \left| \int |u|^2 f dx \right| \leq \epsilon \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + C_\epsilon \|u\|_{L^2(\mathbb{R}^d)}^2, \forall u \in C_0^\infty(\mathbb{R}^d),$$

for all arbitrarily small $\epsilon > 0$ and some $C_\epsilon > 0$. This provides a complete solution to the problem of the infinitesimal form boundedness of the potential energy operator f with respect to the Laplacian $-\Delta$, which is fundamental to quantum mechanics. Its abstract version appears in the so-called KLMN Theorem ([8], Theorem X.17), which is discussed in detail, together with applications to quantum-mechanical Hamiltonian operators and has become an indispensable tool in PDE theory ([7], chap. 5).

Previously, the case of nonnegative f in (1.1) has been studied in a comprehensive way (see e.g. [4], [6], [9], [10]) where different analytic conditions for the so-called trace inequalities of this type can be found.

It is worthwhile to observe that the usual approach is to decompose f into its positive and negative parts : $f = f_+ - f_-$, and to apply the just mentioned results to both f_+ and f_- [6]. However, this procedure drastically diminishes the class of admissible weights f by ignoring a possible cancellation between f_+ and f_- . This cancellation phenomenon is evident for strongly oscillating weights considered below. See for example [11].

One of the main results, we prove that inequality (1.1) is equivalent to the existence of $C > 0$ such that

$$(1.2) \quad |\langle fu, u \rangle| \leq C R^{\frac{2}{1+\beta}} \|\nabla u\|_{L^2(\mathbb{R}^d)}^2, \quad \forall u \in C_0^\infty(B(x_0, R))$$

for all ball $B(x_0, R)$. $B(x_0, R)$ is a Euclidean ball of radius R and centered at x_0 .

Here the "indefinite weight" f may change sign, or even be a complex-valued distribution on \mathbb{R}^d , $d \geq 3$. (In the latter case, the left-hand side (1.1) is understood as $|\langle fu, u \rangle|$, where $\langle f, \cdot \rangle$ is the quadratic form associated with the corresponding multiplication operator f).

We set

$$m_B(g) = \frac{1}{|B(x_0, R)|} \int_{B(x_0, R)} g(y) dy$$

for a ball $B(x_0, R) \subset \mathbb{R}^d$, and denote by $BMO(\mathbb{R}^d)$ the class of $f \in L_{loc}^q(\mathbb{R}^d)$ for which

$$\sup_{R>0} \sup_{x_0 \in \mathbb{R}^d} \frac{1}{|B(x_0, R)|} \int_{B(x_0, R)} |g(y) - m_{B(x_0, R)}(g)|^q dy < +\infty,$$

for any $1 \leq q < \infty$.

Now, we characterize the class of potentials $f \in \mathcal{D}'(\mathbb{R}^d)$ which are there exists $C > 0$ such that (1.2) holds for every ball $B(x_0, R)$.

Theorem 1. *Let $f \in \mathcal{D}'(\mathbb{R}^d)$, $d \geq 2$ and $0 < \beta \leq 1$. Then the following statements are equivalent.*

(1) *There exists a positive constant C such that, for every $\epsilon > 0$,*

$$(1.3) \quad |\langle fu, u \rangle| \leq \epsilon \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + C\epsilon^{-\beta} \|u\|_{L^2(\mathbb{R}^d)}^2, \quad \text{for all } u \in C_0^\infty(\mathbb{R}^d).$$

(2) *There exists a positive constant C such that, for every $R > 0$,*

$$(1.4) \quad |\langle fu, u \rangle| \leq C R^{\frac{2}{1+\beta}} \|\nabla u\|_{L^2(\mathbb{R}^d)}^2, \quad \text{for all } u \in C_0^\infty(B(x_0, R))$$

where C does not depend on x_0 and R .

Define the vector-field $\vec{F} \in \mathcal{D}'(\mathbb{R}^d)^d$ by

$$(1.5) \quad \langle \vec{F}, \vec{\phi} \rangle = - \langle f, \Delta^{-1} \operatorname{div} \vec{\phi} \rangle,$$

for every $\vec{\phi} = (\phi_1, \dots, \phi_d)$ be an arbitrary vector-field in $\mathcal{D} \otimes \mathbb{C}^d$. In particular,

$$(1.6) \quad \langle \vec{F}, \nabla \psi \rangle = - \langle f, \psi \rangle, \quad \psi \in \mathcal{D}(\mathbb{R}^d),$$

i.e.,

$$(1.7) \quad f = \operatorname{div} \vec{F} \text{ in } \mathcal{D}'(\mathbb{R}^d).$$

We have to check that the right-hand side of (1.5) is well-defined, which a priori is not obvious. For $\vec{\phi} \in \mathcal{D} \otimes \mathbb{C}^d$, let

$$w = \Delta^{-1} \operatorname{div} \vec{\phi},$$

where $-\Delta^{-1}g = I_2g$ is the Newtonian potential of $g \in \mathcal{D}$. Clearly,

$$w(x) = O(|x|^{1-d}) \quad \text{and} \quad |\nabla w(x)| = O(|x|^{-d}) \quad \text{as} \quad |x| \rightarrow \infty,$$

and hence

$$w = \Delta^{-1} \operatorname{div} \vec{\phi} = -I_2 \operatorname{div} \vec{\phi} \in \dot{H}^1(\mathbb{R}^d) \cap C^\infty(\mathbb{R}^d).$$

Remark 1. When $f(x) \geq 0$ is locally integrable nonnegative function, Theorem 1 makes it possible to reduce the problem of boundedness for general "indefinite" f to the case of nonnegative weights $|\vec{F}|^2$, which is by now well understood. In particular, combining Theorem 1 and the known criteria in the case $f \geq 0$ (see [4], [6], [9]) we arrive at the following corollary.

Corollary 1. Under the assumptions of Theorem 1, the following statements are equivalent.

(i): Inequality (1.3) holds.

(ii): Suppose that f is represented in the form

$$(1.8) \quad f = \operatorname{div} \vec{F},$$

where $\vec{F} = \nabla \Delta^{-1} f \in L^2_{loc}(\mathbb{R}^d)^d$ and the measure $\mu \in \mathcal{M}^+(\mathbb{R}^d)$ defined by

$$(1.9) \quad d\mu = |\vec{F}(x)|^2 dx$$

has the property that, there exists $C > 0$ such that

$$(1.10) \quad \left| \int |u(x)|^2 d\mu \right| \leq \epsilon \|\nabla u\|_{L^2(\mathbb{R}^d)}^2 + C\epsilon^{-\beta} \|u\|_{L^2(\mathbb{R}^d)}^2, \quad \forall u \in C_0^\infty(\mathbb{R}^d)$$

for every $\epsilon > 0$.

(iii): For μ defined by (1.9),

$$\lim_{R \rightarrow 0^+} \sup_{x_0 \in \mathbb{R}^d} \left\{ \frac{\|\mu_{B(x_0, R)}\|_{\dot{H}^{-1}(\mathbb{R}^d)}^2}{\mu(B(x_0, R))} \right\} = 0,$$

where $\mu_{B(x_0, R)}$ is the restriction of μ to the ball $B(x_0, R)$.

Before proceeding to our main result, it is instructive to demonstrate the cancellation phenomenon mentioned above by considering an example of a strongly oscillating weight.

Example 1. Let us set

$$f(x) = |x|^{d-2} \sin(|x|^d),$$

where $d \geq 3$ is an integer, which may be arbitrary large. Obviously, both f_+ and f_- fail to satisfy (1.3) due to the growth of the amplitude at infinity. However,

$$(1.11) \quad f(x) = \operatorname{div} \vec{F}(x) + O(|x|^{-2}), \quad \text{where} \quad \vec{F}(x) = -\frac{1}{d} \frac{\vec{x}}{|x|^2} \cos(|x|^d).$$

By Hardy's inequality in \mathbb{R}^d , $d \geq 3$,

$$\int_{\mathbb{R}^d} \frac{|u(x)|^2}{|x|^2} dx \leq \frac{4}{(d-2)^2} \|\nabla u\|_{L^2}^2, \quad u \in C_0^\infty(\mathbb{R}^d),$$

and hence the term $O(|x|^{-2})$ in (1.11) is harmless, whereas \vec{F} clearly satisfies (1.10) since $|\vec{F}(x)|^2 \leq |x|^{-2}$. This shows that f is admissible for (1.2), while $|f|$ is obviously not (see [6]).

Theorem 2. Let f be a complex-valued distribution on \mathbb{R}^d , $d \geq 3$ and let $0 < \beta \leq 1$. Then (1.4) holds if and only if f is the divergence of a vector-field $\vec{F} : \mathbb{R}^d \rightarrow \mathbb{C}^d$ such that

$$(1.12) \quad \int_{B(x_0, R)} \left| \vec{F}(x) - m_{B(x_0, R)}(\vec{F}) \right|^2 dx \leq \text{const } R^{d-2+\frac{4}{1+\beta}}, \quad \text{for all } R > 0.$$

where the constant is independent of x_0 and R . The vector-field $\vec{F} \in L_{loc}^2(\mathbb{R}^d)^d$ can be chosen as $\vec{F} = \nabla \Delta^{-1} f$ (see [6]).

Remark 2.

1.: In case, $\beta = 1$, it follows that (1.12) holds if and only if $\vec{F} \in BMO(\mathbb{R}^d)^d$.

In order, $f \in BMO^{-1}(\mathbb{R}^d) = \dot{F}_{\infty}^{-1,2}(\mathbb{R}^d)$, where $\dot{F}_q^{r,p}$ stands for the scale of homogeneous Triebel-Lizorkin spaces (see [13]). Similarly, in the case $0 < \beta < 1$, (1.12) holds if and only if \vec{F} is Hölder-continuous :

$$\left| \vec{F}(x) - \vec{F}(y) \right| \leq c|x - y|^{\frac{1-\beta}{1+\beta}}, \quad |x - y| < R.$$

2.: In the case $\beta = 1$, statement (i) of Theorem 2 (sufficiency of the condition $\vec{F} \in BMO(\mathbb{R}^d)^d$) is equivalent via the $\mathcal{H}^1 - BMO$ duality to the inequality

$$\|u \nabla u\|_{\mathcal{H}^1(\mathbb{R}^d)} \leq C \|u\|_{L^2(\mathbb{R}^d)} \|\nabla u\|_{L^2(\mathbb{R}^d)}, \quad \forall u \in C_0^\infty(\mathbb{R}^d).$$

Here $\mathcal{H}^1(\mathbb{R}^d)$ is the real Hardy space on \mathbb{R}^d [12]. The preceding estimate yields the following vector-valued inequality which is used in studies of the Navier-Stokes equation, and is related to the compensated compactness phenomenon (see [1]) :

$$\begin{aligned} \|(\vec{u} \cdot \nabla) \vec{u}\|_{\mathcal{H}^1(\mathbb{R}^d)} &\leq C \|\vec{u}\|_{L^2(\mathbb{R}^d)^d} \|\nabla \vec{u}\|_{L^2(\mathbb{R}^d)^d} \\ \operatorname{div} \vec{u} &= \vec{0}, \quad \forall \vec{u} \in C_0^\infty(\mathbb{R}^d)^d. \end{aligned}$$

Before proving the theorem, let us established certain localized versions of the necessary condition for (1.4). Set

$$\omega_{R,x_0}(x) = \omega\left(\frac{x - x_0}{R}\right)$$

where $\omega \in C_0^\infty(B(0, 1))$ is a smooth cut-off function with the following properties

$$|\omega(x)| \leq 1 \text{ and } |\nabla \omega(x)| \leq 1 \text{ for } x \in B(0, 1).$$

With this definition, we obtain the following more useful statement.

Proposition 1. Suppose $f \in \mathcal{D}'(\mathbb{R}^d)$ and $0 < \beta \leq 1$. Suppose that (1.4) holds for every $R \in (0, +\infty)$. Let \vec{F} be defined by $\vec{F} = \nabla \Delta^{-1} f$.

(a): For $d \geq 3$,

$$\int_{\mathbb{R}^d} |\nabla \Delta^{-1}(\omega_{R,x_0} f)|^2 dx \leq CR^{d-2+\frac{4}{1+\beta}}, \quad 0 < R < +\infty$$

(b): For $d \geq 2$,

$$\int_{B(x_0, R)} |\nabla \Delta^{-1} (\omega_{R, x_0} f)|^2 dx \leq C R^{d-2+\frac{4}{1+\beta}}, \quad 0 < R < +\infty$$

Now we can state the following

Lemma 1. Suppose $f \in \mathcal{D}'(\mathbb{R}^d)$, $d \geq 2$ and $0 < \beta \leq 1$. Suppose that (1.4) holds for every $R \in (0, +\infty)$. Then we have

$$\int_{B(x_0, R)} |\nabla \Delta^{-1} f - m_{B(x_0, R)}(\nabla \Delta^{-1} f)|^2 dx \leq C R^{d-2+\frac{4}{1+\beta}}.$$

We are now in a position to give the proof of theorem 2. We need only to prove the statement (i) since (ii) follow from Proposition 1 and Lemma 1.

PROOF. Suppose that f is represented in the form (1.7) so that (1.12) is satisfied for all $R > 0$. Applying the multiplicative inequality nonnegative measures ([5], th.1.4.7) to $|\vec{F}|^2 dx$, we get :

$$\int_{B(x_0, R)} |\vec{F}(x)|^2 |u(x)|^2 dx \leq C \|\nabla u\|_{L^2(\mathbb{R}^d)}^{2(\frac{\beta-1}{\beta+1})} \|u\|_{L^2(\mathbb{R}^d)}^{\frac{4}{\beta+1}}.$$

Hence,

$$\begin{aligned} |\langle fu, u \rangle| &= |\langle \vec{F}u, \nabla u \rangle| \leq \|\vec{F}u\|_{L^2(\mathbb{R}^d)} \|\nabla u\|_{L^2(\mathbb{R}^d)} \\ &\leq C_1^{\frac{1}{2}} \|\nabla u\|_{L^2(\mathbb{R}^d)}^{1+\frac{\beta-1}{\beta+1}} \|u\|_{L^2(\mathbb{R}^d)}^{\frac{2}{\beta+1}} \end{aligned}$$

Combining the preceding estimates with the following inequality ([7], th 3.2.1) :

$$\|u\|_{L^2} \leq C(d)R \|\nabla u\|_{L^2}, \quad u \in C_0^\infty(B(x_0, R)),$$

we get

$$|\langle fu, u \rangle| \leq C R^{\frac{2}{1+\beta}} \|\nabla u\|_{L^2}^2, \quad u \in C_0^\infty(B(x_0, R)).$$

The proof of theorem 2 is complete. \square

We use know characterizations of the Morrey-Campanato spaces. In particular,

Proposition 2. For $0 < \beta < 1$, condition (1.12) is equivalent to the condition $\vec{F} \in \Lambda_\gamma(\mathbb{R}^d)$ where $\gamma = \frac{1-\beta}{1+\beta}$. In the case $\beta = 1$, we have $\vec{F} \in BMO(\mathbb{R}^d)^d$.

It is easy to see that in the case $\beta = 1$, the sufficiently part of Theorem 2 is equivalent to inequality :

$$\left| \langle \vec{F}u, \nabla u \rangle \right| \leq C \left\| \vec{F} \right\|_{BMO(\mathbb{R}^d)^d} \|u\|_{L^2(\mathbb{R}^d)} \|\nabla u\|_{L^2(\mathbb{R}^d)}, \forall u \in C_0^\infty(\mathbb{R}^d)$$

By duality, the preceding inequality yields :

$$\|u \nabla u\|_{\mathcal{H}^1(\mathbb{R}^d)} \leq C \|u\|_{L^2(\mathbb{R}^d)} \|\nabla u\|_{L^2(\mathbb{R}^d)}, \quad \forall u \in C_0^\infty(\mathbb{R}^d).$$

where $\mathcal{H}^1(\mathbb{R}^d)$ is a real Hardy space [12]. Such inequalities are useful in hydrodynamics [1]. As an immediate consequence, we obtain the vector-valued quadratic form :

$$\begin{aligned} \|(\vec{u} \cdot \nabla) \vec{u}\|_{\mathcal{H}^1(\mathbb{R}^d)} &\leq C \|\vec{u}\|_{L^2(\mathbb{R}^d)^d} \|\nabla \vec{u}\|_{L^2(\mathbb{R}^d)^d} \\ \operatorname{div} \vec{u} &= \vec{0}, \quad \forall \vec{u} \in C_0^\infty(\mathbb{R}^d)^d. \end{aligned}$$

Both of the preceding inequalities are corollaries of the homogeneous version of the "div - curl" Lemma [1]. The following corollary which is an immediate consequence of Theorem 2 and the characterizations of Morrey-Campanato spaces [3], gives a necessary and sufficient condition for (1.12) in terms of homogeneous Besov spaces of negative order.

Corollary 2. *Under the assumptions of Theorem 2, in the case $\beta = 1$, condition (1.12) is equivalent to $f \in BMO^{-1}(\mathbb{R}^d)$. Similarly, in the case $0 < \beta < 1$, condition (1.12) is equivalent to $f \in \dot{B}_{\infty}^{-\frac{2\beta}{1+\beta}, \infty}(\mathbb{R}^d)$.*

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