Inverse scattering for Schrödinger equation in the frequency domain via data-driven reduced order modeling

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Motivation: inverse scattering



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- Reduced order model (ROM) framework for inverse scattering:
- (1) Construct a **data-driven ROM** from the data
- (2) Formulate inverse scattering problem as **ROM misfit** minimization
 - Similar approach works very well in the time domain setting
 - Extend to frequency domain



• Classical inverse scattering: Schrödinger equation

$$\left[-\Delta+q(\mathbf{x})-k^2
ight]u(\mathbf{x};k)=0,\quad\mathbf{x}\in\mathbb{R}^d,\quad d=2,3$$

Wavefield decomposition: incoming + scattered

$$u(\mathbf{x}; k) = u^{\text{inc}}(\mathbf{x}; k) + u^{\text{scat}}(\mathbf{x}; k)$$

• Radiation condition, e.g., Sommerfeld

$$\lim_{r\to\infty}r^{\frac{d-1}{2}}\left(\frac{\partial}{\partial r}-\imath k\right)u^{\mathrm{scat}}(\mathbf{x};k)=0,$$

• Measure $u^{\text{scat}}(\mathbf{x}; k)$ at $r = ||\mathbf{x}|| \to \infty$ for various $u^{\text{inc}}(\mathbf{x}; k)$ and k to estimate the scattering potential $q(\mathbf{x})$

Our formulation: forward model

• Schrödinger equation in Ω only:

$$\left[-\Delta+q(\mathbf{x})-k^2
ight]u^{(s)}(\mathbf{x};k)=0,\quad\mathbf{x}\in\Omega,$$

 The radiation condition is approximated by a first order absorbing boundary condition + sources

$$[\mathbf{n}(\mathbf{x})\cdot\nabla - \imath k] u^{(s)}(\mathbf{x};k) = p_s(\mathbf{x}), \ \mathbf{x} \in \partial\Omega, \ s = 1, \dots, m,$$

where the sources

$$p_{\boldsymbol{s}}(\boldsymbol{x}) = \xi^{(\boldsymbol{s})}(\boldsymbol{x}) \left[\boldsymbol{\mathsf{n}}(\boldsymbol{x}) \cdot \nabla - \imath k \right] u^{\mathsf{inc},(\boldsymbol{s})}(\boldsymbol{x};k), \ \boldsymbol{x} \in \partial \Omega, \ \boldsymbol{s} = 1, \dots, m,$$

are restrictions of the incoming waves to "windows" on $\partial \Omega$



• For *n* sampling wavenumbers

$$0 < k_1 < k_2 < \ldots < k_n$$

illuminate the medium with each of *m* sources $p_s(\mathbf{x})$, s = 1, ..., m, and measure the **boundary data**

$$\mathcal{D} = \left\{ \phi_j^{(s)} = \left. u_j^{(s)} \right|_{\partial \Omega}, \quad \partial_k \phi_j^{(s)} = \left. \partial_k u_j^{(s)} \right|_{\partial \Omega} \right\}_{j=1,\dots,n; \ s=1,\dots,m}$$

 Inverse Scattering Problem (ISP): estimate the scattering potential q(x) in Ω from the data D



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• Consider wavefield snapshots for all frequencies and sources

$$u_{j}^{(s)}(\mathbf{x}) = u^{(s)}(\mathbf{x}; k_{j}), \quad j = 1, \dots, n, \quad s = 1, \dots, m$$

• We want to get a ROM of the forward problem by projecting onto

$$\mathcal{X} = \operatorname{span}\left\{u_{j}^{(s)}\right\}_{j=1,\dots,n;\ s=1,\dots,m}$$

 The ROM consists of mn × mn mass (M), stiffness (S) and boundary (B) matrices with blocks

$$\begin{split} [\mathbf{s}_{ij}]_{rs} &= \int_{\Omega} \overline{\nabla u_i^{(r)}} \cdot \nabla u_j^{(s)} d\mathbf{x} + \int_{\Omega} q \overline{u_i^{(r)}} u_j^{(s)} d\mathbf{x}, \\ [\mathbf{m}_{ij}]_{rs} &= \int_{\Omega} \overline{u_i^{(r)}} u_j^{(s)} d\mathbf{x}, \quad [\mathbf{b}_{ij}]_{rs} = \int_{\partial \Omega} \overline{u_i^{(r)}} u_j^{(s)} d\Sigma, \end{split}$$

and takes the form

$$(\mathbf{S} - k^2 \mathbf{M} + \imath k \mathbf{B}) \widetilde{\mathbf{u}}^{(s)} = \mathbf{p}^s, \quad s = 1, \dots, m$$



Data-driven ROM

- Big issue: wavefield snapshots are unknown in Ω, only measured at ∂Ω. How to compute the ROM?
- Need a data-driven process for ROM computation
- Two-stage process:

First, given the data \mathcal{D} , compute **boundary integrals**

$$\begin{split} [\mathbf{b}_{ij}]_{rs} &= \int_{\partial\Omega} \overline{\phi_i^{(r)}} \phi_j^{(s)} d\Sigma, \\ [\mathbf{c}_j]_{rs} &= \int_{\partial\Omega} \left[-\overline{\phi_j^{(r)}} \partial_k \phi_j^{(s)} + \phi_j^{(s)} \overline{\partial_k \phi_j^{(r)}} \right] d\Sigma, \\ [\mathbf{d}_j]_{rs} &= \int_{\partial\Omega} p_r \phi_j^{(s)} d\Sigma, \\ [\partial_k \mathbf{d}_j]_{rs} &= \int_{\partial\Omega} p_r \partial_k \phi_j^{(s)} d\Sigma, \end{split}$$

for i, j = 1, ..., n, r, s = 1, ..., m

Data-driven ROM

• Second stage: from the boundary integrals compute the blocks of stiffness and mass matrices:

$$\begin{split} \mathbf{s}_{ij} &= \frac{k_i^2 \mathbf{d}_i^* - k_j^2 \mathbf{d}_j}{k_i^2 - k_j^2} - i \frac{(k_i k_j^2 + k_i^2 k_j) \mathbf{b}_{ij}}{k_i^2 - k_j^2}, \quad i \neq j \\ \mathbf{s}_{jj} &= \frac{1}{2} \Big(k_j \Re(\partial_k \mathbf{d}_j) + 2 \Re(\mathbf{d}_j) \Big) + \frac{i k_j^2}{2} \mathbf{c}_j, \\ \mathbf{m}_{ij} &= \frac{\mathbf{d}_i^* - \mathbf{d}_j}{k_i^2 - k_j^2} - i \frac{\mathbf{b}_{ij}}{k_j - k_i}, \quad i \neq j \\ \mathbf{m}_{jj} &= \frac{1}{2k_j} \Re(\partial_k \mathbf{d}_j) + \frac{i}{2} \mathbf{c}_j, \end{split}$$

for i, j = 1, ..., n.

- Resemble classical Löwner product formulas
- Note that b_{ij} and c_j are non-linear in the data
- Derivatives w.r.t. k only needed for digonal blocks



Recall the blocks of the stiffness matrix

$$[\mathbf{s}_{ij}]_{rs} = \int_{\Omega} \overline{\nabla u_i^{(r)}} \cdot \nabla u_j^{(s)} d\mathbf{x} + \int_{\Omega} q \overline{u_i^{(r)}} u_j^{(s)} d\mathbf{x}$$

- This is not an orthogonal projection ROM, since the snapshots u_j are not orthonormal
- Can transform the ROM to an orthogonal projection ROM via block Lanczos process applied to S with M-inner product to obtain

$$[\mathbf{t}_{ij}]_{rs} = \int_{\Omega} \overline{\nabla v_i^{(r)}} \cdot \nabla v_j^{(s)} d\mathbf{x} + \int_{\Omega} q \overline{v_i^{(r)}} v_j^{(s)} d\mathbf{x},$$

the blocks of a **block-tridiagonal** matrix **T**, where $\{v_j\}_{j=1}^n$ is an orthonormal basis for the projection space \mathcal{X}

Using ROM to solve ISP

- How can the ROM be used to solve the ISP?
- Conventional data misfit a.k.a. Full Waveform Inversion (FWI) is a non-linear least squares problem

$$\underset{\widehat{q} \in \mathcal{Q}}{\text{minimize}} \sum_{j=1}^{n} \sum_{s=1}^{m} \int_{\partial \Omega} \left| \phi_{j}^{(s),\text{meas}} - \phi_{j}^{(s)}[\widehat{q}] \right|^{2} d\Sigma$$

• We propose instead to use

$$\min_{\widehat{\boldsymbol{q}} \in \mathcal{Q}} \mathcal{F}(\widehat{\boldsymbol{q}}) + \mu \mathcal{R}(\widehat{\boldsymbol{q}}),$$

with either

$$\mathcal{F}_{\mathsf{S}}(\widehat{q}) = \left\| \mathsf{Triu}(\mathsf{S}^{\mathsf{meas}} - \mathsf{S}[\widehat{q}]) \right\|_{2}^{2}$$

or

$$\mathcal{F}_{\mathsf{T}}(\widehat{q}) = \left\|\mathsf{Triu}(\mathsf{T}^{\mathsf{meas}} - \mathsf{T}[\widehat{q}])\right\|_2^2$$



Advantages of using ROMs for solving ISP

- Why is ROM-based ISP better than conventional frequency domain FWI?
- Data least squares misfit functional is known to be highly non-convex, leading to optimization getting stuck in local minima, slow convergence, dependence on initial guess and other issues
- ROM misfit is much better behaved in the time domain case
- We expect similar behavior in the frequency domain case



- $q(\mathbf{x})$ with two inclusions in the unit square $\Omega = [0, 1] \times [0, 1]$
- *m* = 8 sources along the top boundary of Ω
- *n* = 8 sampling wavenumbers *k* = 20, 25, 30, 35, 40, 45, 50, 55
- 2.5% noise added to the data, ROM computation regularized with **spectral truncation**
- Search space Q with 20 \times 20 = 400 basis functions
- 10 Gauss-Newton iterations regularized with Tikhonov with adaptive choice of μ



Numerical example: results





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Numerical example: quality control



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Conclusions and future work

- We introduced ROM framework numerical solution of inverse scattering for Schrödinger equation with frequency domain measurements
- Separate the ISP into two steps:
 - Construct the ROM from data
 - Use ROM misfit as optimization objective
- Much better numerical results than conventional FWI least squares data misfit

Future work:

- Extend to Helmholtz equation to account for kinematic effects
- Use the ROM to estimate internal solutions to use in a modified Lippmann-Schwinger formulation for solving the IPS

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